

# **ELEN E3106/4106 Lecture 8**

## **p-n Junctions Part I**

### **Outline**

- Simple theory to p-n junctions
- Junctions in equilibrium
- Contact Potential
- Depletion Approximation
- Electrostatics Calculations & Diagrams

### **Assignments:**

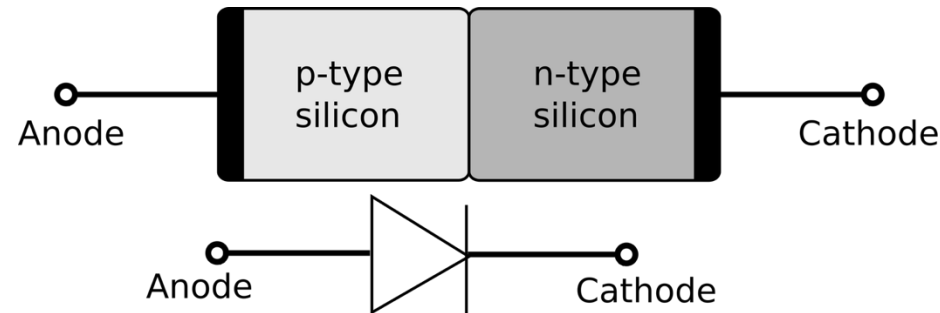
Reading: Streetman and Banerjee §5.2-5.3

Homework 3 due tomorrow, Friday Sep 26<sup>th</sup> by 5pm

Exam 1 this Tuesday Sept. 30<sup>th</sup>

# p-n Junction

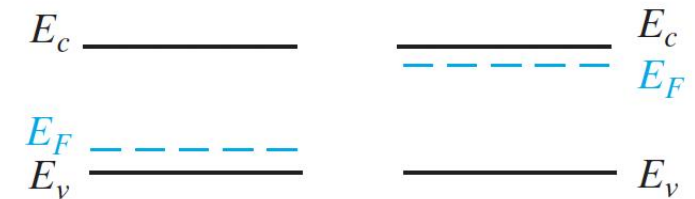
- Most semiconductor devices contain at least one junction between \_\_\_\_\_ and \_\_\_\_\_ material
- So what happens when we bring together a slab of n-type material and a slab of p-type material?
- We already have most of the tools to understand this
- An understanding of p-n junctions is needed to analyze the other devices we will look at in this class



Assumptions for this lecture:

# Equilibrium Conditions: Isolated Regions

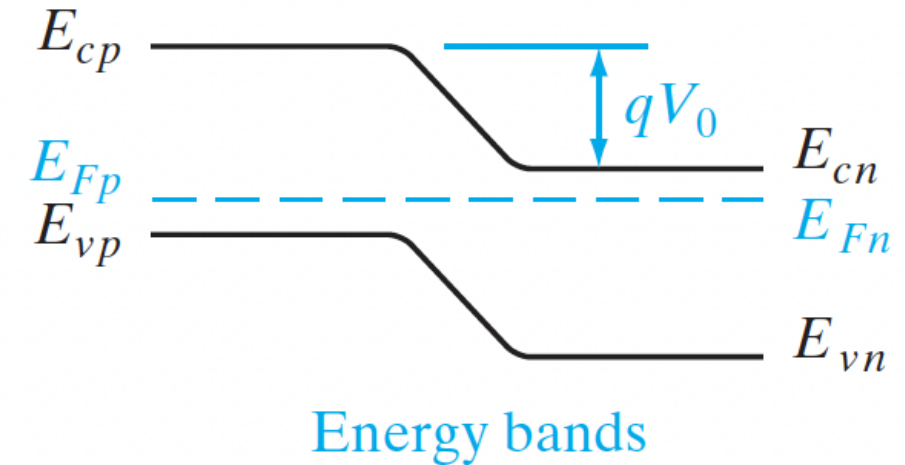
- Recall: If you have isolated regions of n- and p-type material, we know what our band diagrams look like
- Equilibrium means no \_\_\_\_\_ and no \_\_\_\_\_
- First, we will look at the \_\_\_\_\_ *junction*
  - \_\_\_\_\_ p doping on one side of a sharp junction and uniform n doping on the other side
- Before joining, the n-type material has a large concentration of \_\_\_\_\_
- The p-type material has a large concentration of \_\_\_\_\_



(a)

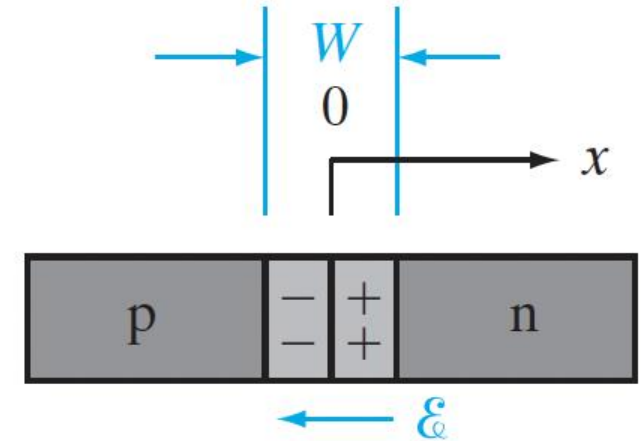
# Equilibrium Conditions: Step/Abrupt Junction

- Immediately upon joining, there is a large \_\_\_\_\_ of carriers across the sample
- Holes \_\_\_\_\_ from the p side into the n side
- Electrons \_\_\_\_\_ from n to p
- An opposing \_\_\_\_\_ is created at the junction and builds up until the net current is \_\_\_\_\_
  - Recall:
$$J_p(\text{drift}) + J_p(\text{diff.}) = 0$$
$$J_n(\text{drift}) + J_n(\text{diff.}) = 0$$
- What does the Fermi level look like in equilibrium?
- We can now draw the band diagram



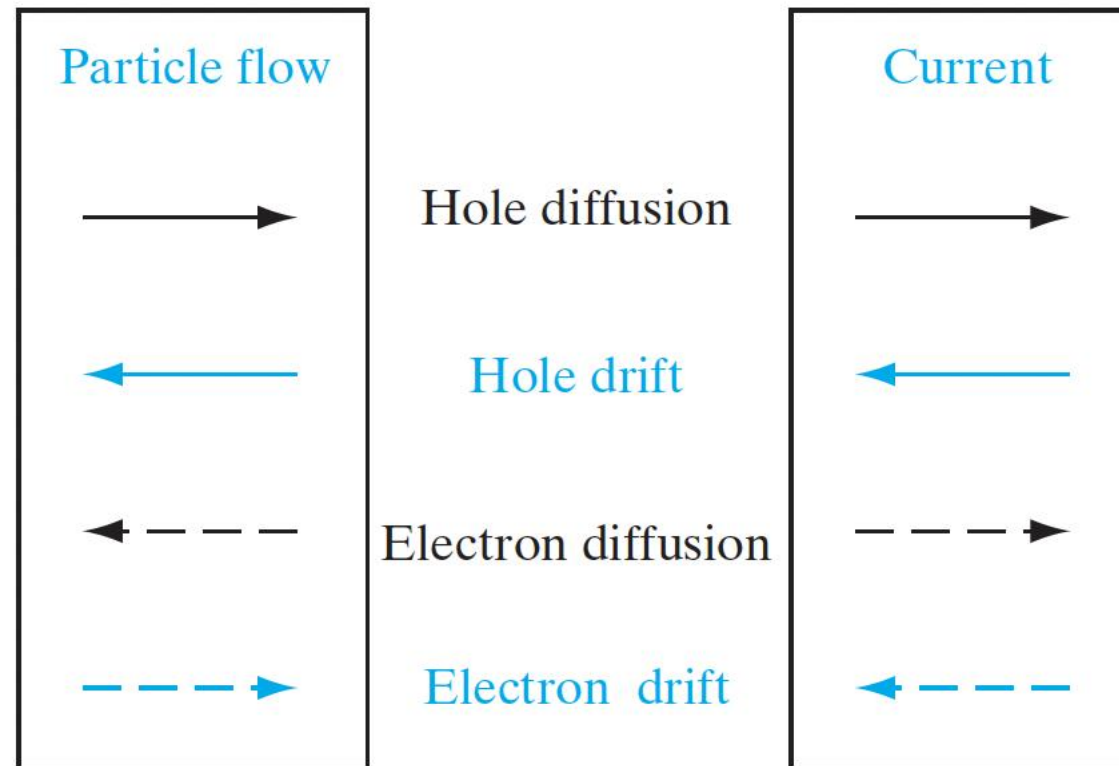
# Space Charge in the Transition Region

- e<sup>-</sup> diffusing from n to p side will leave behind \_\_\_\_\_
- h<sup>+</sup> diffusing from p to n side will leave behind \_\_\_\_\_
- After diffusion is done,
- This leads to the development of a region of positive \_\_\_\_\_ near the n side of the junction and negative charge near the p side
- This region is denoted \_\_\_\_\_, also called the \_\_\_\_\_
- The E-field appears across this region near the \_\_\_\_\_



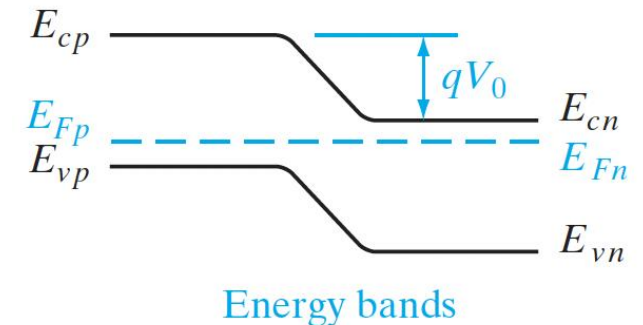
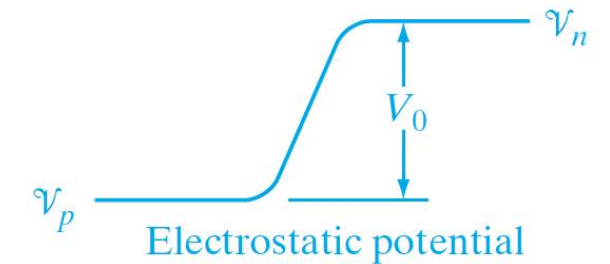
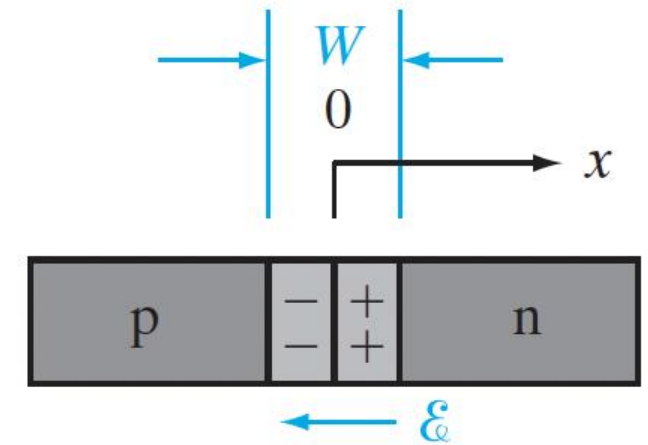
# Visualizing Directions of Particle and Current Flow

- Directions of the four components of particle flow within the transition region, and the resulting current directions
- Recall: electron \_\_\_\_\_ is opposite to the \_\_\_\_\_ of flow of electrons



# Contact Potential

- There is an equilibrium potential difference \_\_\_\_\_ across  $W$
- Gradient in potential oppose to the direction of E-field
- We assume E-field is \_\_\_\_\_ in the \_\_\_\_\_ regions outside  $W$
- The contact potential, \_\_\_\_\_ =  $V_n - V_p$ 
  - $V_0$  is a built-in potential barrier
  - Necessary to maintain equilibrium at the junction
- Can you measure it with a voltmeter?



# Calculating Contact Potential

- From the band diagram,

$$qV_0 = E_{vp} - E_{vn}$$

$$qV_0 = (E_{ip} - E_F) + (E_F - E_{in})$$

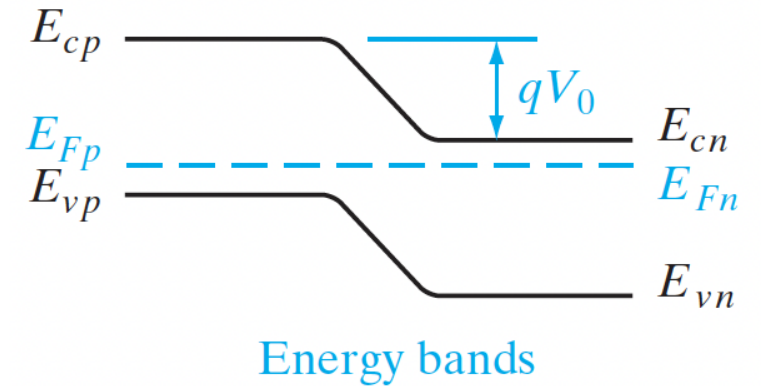
- We can write  $V_0$  in terms of the the acceptor and donor concentrations on the p- and n-sides, respectively:

$$V_0 = \frac{kT}{q} \ln \frac{N_a}{n_i^2 / N_d} = \frac{kT}{q} \ln \frac{N_a N_d}{n_i^2}$$

- Equilibrium conditions still hold true:

$$p_p n_p = n_i^2 = p_n n_n$$

- What do the subscripts in n and p denote?



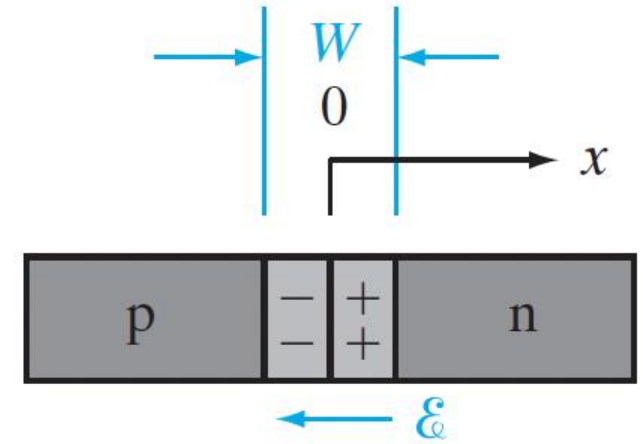
$$E_{ip} - E_F = kT \ln \frac{p_p}{n_i}$$

$$E_F - E_{in} = kT \ln \frac{n_n}{n_i}$$



## Carriers in a p-n Junction

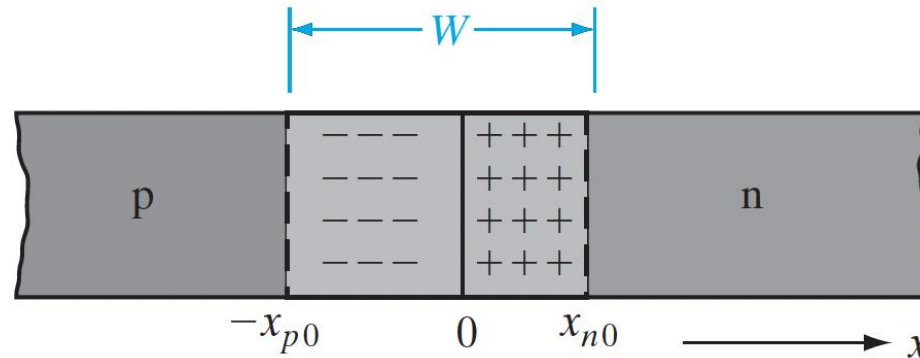
- Notation:
  - On the p-side majority carrier:  $p_p = \underline{\hspace{1cm}}$  outside of W
  - On the n-side, majority carrier  $n_n = \underline{\hspace{1cm}}$  outside of W
  
- Using  $n_p p_p = n_i^2$  on p-side, minority carriers there
  - $n_p =$
  
- From the built-in voltage:  $\frac{p_p}{p_n} = \frac{n_n}{n_p} = e^{qV_0 / kT}$
  
- This relates the majority/minority carrier concentration on either side of the junction
  
- Which becomes more useful next lecture(s) when we apply an \_\_\_\_\_



# Depletion Approximation

- We approximate there is an \_\_\_\_\_ between the space charge ( $N_D - N_A$ ) region and the two quasi-neutral (n and p) regions
- $e^-$  and  $h^+$  are in transit in  $W$  from one side of the junction to the other. Some electrons diffuse from n to p
- The \_\_\_\_\_ serves to sweep out carriers which have wandered into  $W$
- Results: There are \_\_\_\_\_ in the transition region  $W$  at any time
- We consider the space charge within  $W$  is due only to \_\_\_\_\_ and \_\_\_\_\_

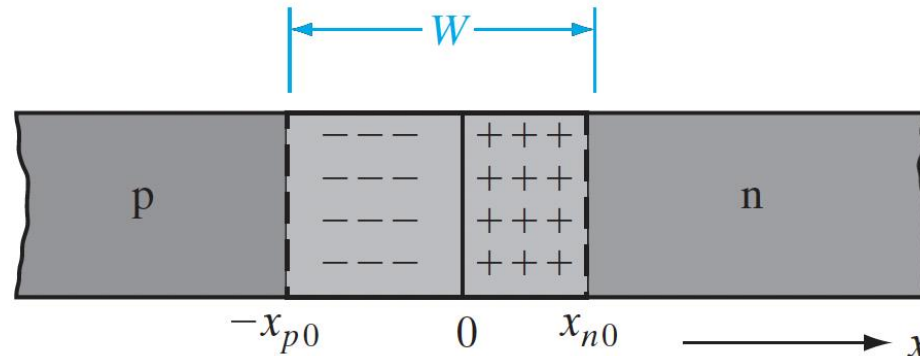
***Space charge and  $E$ -field distribution within the transition region of a p-n junction with  $N_d > N_a$ :***



# Depletion Approximation Continued

- What is the depletion region?
  - What is the space charge region (SCR)?
  - What are the quasi-neutral regions (QNR)?
- 
- If the SCR width is  $W = x_p + x_n$ , do the two ( $x_p$ ,  $x_n$ ) sides have to be equal? Why or why not?

***Space charge and E-field distribution within the transition region of a p-n junction with  $N_d > N_a$ :***

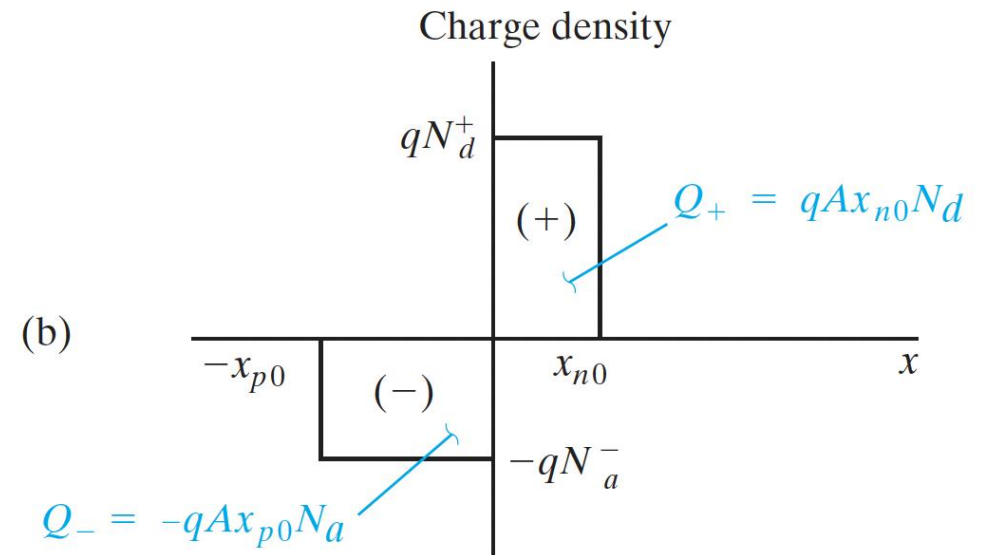
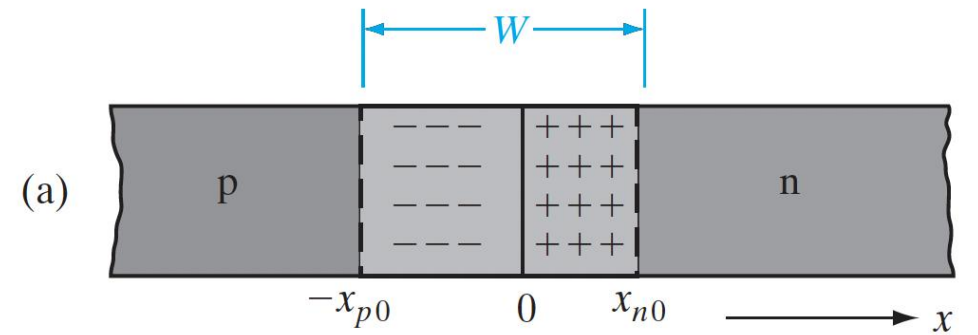


# p-n Junction Electrostatics: Charge

Given a sample with cross-sectional area  $A$ ,

- Charge on the p-side:
- Charge on the n-side:
- Total charge on either side of the junction:

$$qAx_{p0}N_a = qAx_{n0}N_d$$



# p-n Junction Electrostatics: E-Field

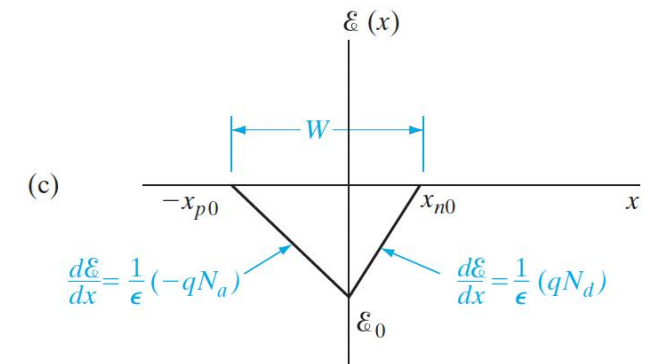
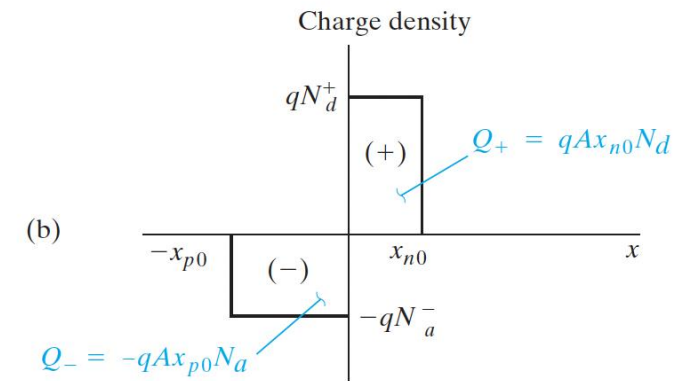
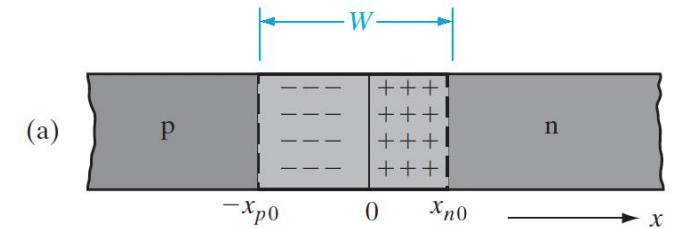
- To calculate the E-field within  $W$ , we can use Poisson's equation, which relates the gradient of the E-field to the local space charge at any point  $x$ :

$$\frac{d\mathcal{E}(x)}{dx} = \frac{q}{\epsilon}(p - n + N_d^+ - N_a^-)$$

- If we neglect  $p$  and  $n$  in the SCR, and assuming complete ionization, we can simplify

$$\frac{d\mathcal{E}}{dx} = \frac{q}{\epsilon}N_d, \quad 0 < x < x_{n0}$$

$$\frac{d\mathcal{E}}{dx} = -\frac{q}{\epsilon}N_a, \quad -x_{p0} < x < 0$$



# p-n Junction Electrostatics: E-Field Continued

- We can easily relate E-field to \_\_\_\_\_ (the field at any x is the negative of the potential gradient at that point)

$$\mathcal{E}(x) = -\frac{dV(x)}{dx} \quad \text{or} \quad -V_0 = \int_{-x_{p0}}^{x_{n0}} \mathcal{E}(x) dx$$

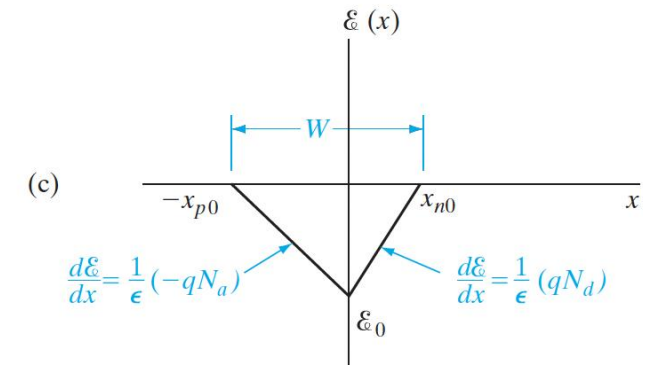
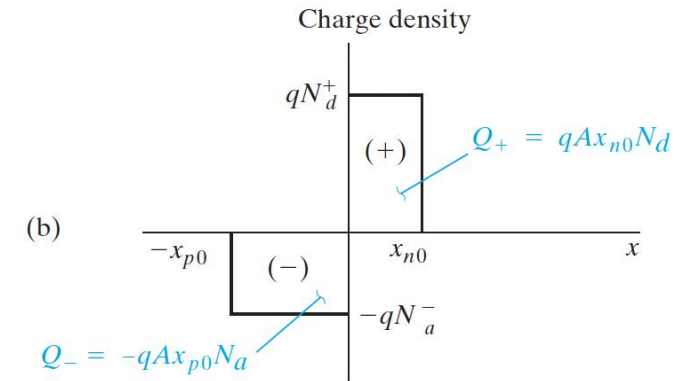
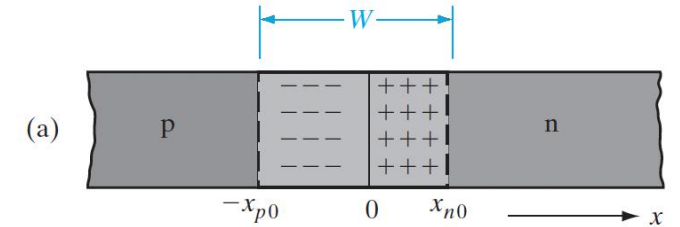
- Therefore, the negative of the contact potential is simply the \_\_\_\_\_ under the E(x) versus x triangle!

- This relates the contact potential to the width of the depletion region:

$$V_0 = -\frac{1}{2} \mathcal{E}_0 W = \frac{1}{2} \frac{q}{\epsilon} N_d x_{n0} W$$

- Which can also be written,

$$V_0 = \frac{1}{2} \frac{q}{\epsilon} \frac{N_a N_d}{N_a + N_d} W^2$$



# p-n Junction Electrostatics: E-Field

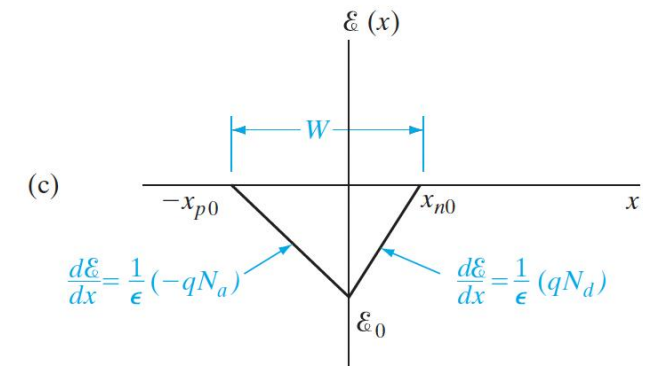
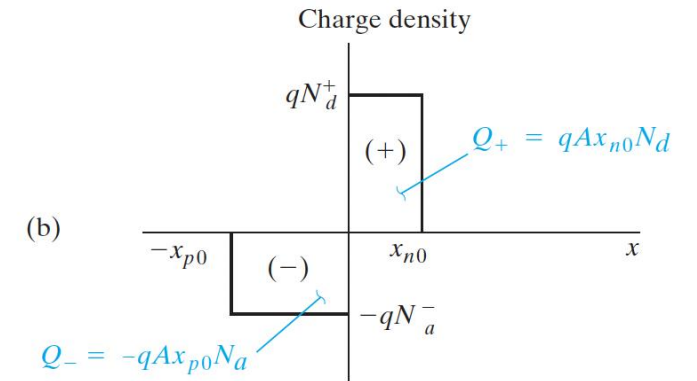
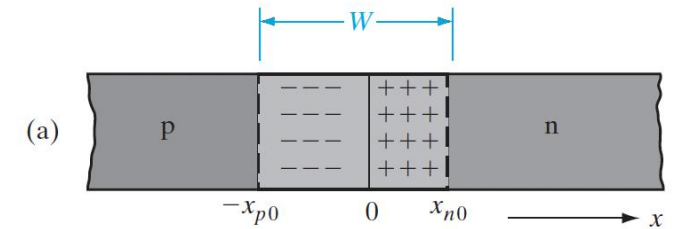
- What does the E-field look like?
- \_\_\_\_\_ slope (E-field increasing with x) on the n-side
- \_\_\_\_\_ slope (E-field becomes more negative as x increases) on the p-side
- Max value  $E_0$  at  $x = 0$

$$\int_{\mathcal{E}_0}^0 d\mathcal{E} = \frac{q}{\epsilon} N_d \int_0^{x_{n0}} dx, \quad 0 < x < x_{n0}$$

$$\mathcal{E}_0 = -\frac{q}{\epsilon} N_d x_{n0} = -\frac{q}{\epsilon} N_a x_{p0}$$

$$\int_0^{\mathcal{E}_0} d\mathcal{E} = -\frac{q}{\epsilon} N_a \int_{-x_{p0}}^0 dx, \quad -x_{p0} < x < 0$$

- We know E-field points in  $-x$  direction (from n to p), so it is \_\_\_\_\_
- E-field assumed to go to zero outside of W!



# p-n Junction Electrostatics: Depletion Widths

- What about the depletion width,  $W$ ?
- Rearranging  $V_0$  equation,

$$W = \left[ \frac{2\epsilon V_0}{q} \left( \frac{N_a + N_d}{N_a N_d} \right) \right]^{1/2} = \left[ \frac{2\epsilon V_0}{q} \left( \frac{1}{N_a} + \frac{1}{N_d} \right) \right]^{1/2} \quad W = \left[ \frac{2\epsilon kT}{q^2} \left( \ln \frac{N_a N_d}{n_i^2} \right) \left( \frac{1}{N_a} + \frac{1}{N_d} \right) \right]^{1/2}$$

- Note: transition/depletion width  $W$  varies as the \_\_\_\_\_
- We can also calculate the penetration of the transition region into the n and p materials:

$$x_{p0} = \frac{W N_d}{N_a + N_d} = \frac{W}{1 + N_a/N_d} = \left\{ \frac{2\epsilon V_0}{q} \left[ \frac{N_d}{N_a(N_a + N_d)} \right] \right\}^{1/2}$$

$$x_{n0} = \frac{W N_a}{N_a + N_d} = \frac{W}{1 + N_d/N_a} = \left\{ \frac{2\epsilon V_0}{q} \left[ \frac{N_a}{N_d(N_a + N_d)} \right] \right\}^{1/2}$$

- As we expect, this predicts transition region extends further into side with \_\_\_\_\_! (i.e. if  $N_a \ll N_d$ ,  $x_{p0}$  is large compared with  $x_{n0}$ )



# Problem: p-n Junctions in Equilibrium

- An abrupt Si p-n junction has  $N_a = 10^{18} \text{ cm}^{-3}$  on one side and  $N_d = 5 \times 10^{15} \text{ cm}^{-3}$  on the other. Calculate the Fermi level positions at 300 K in the p and n regions. Draw an equilibrium band diagram for the junction and determine the contact potential  $V_0$  from the diagram. Compare the results with  $V_0$  calculation.

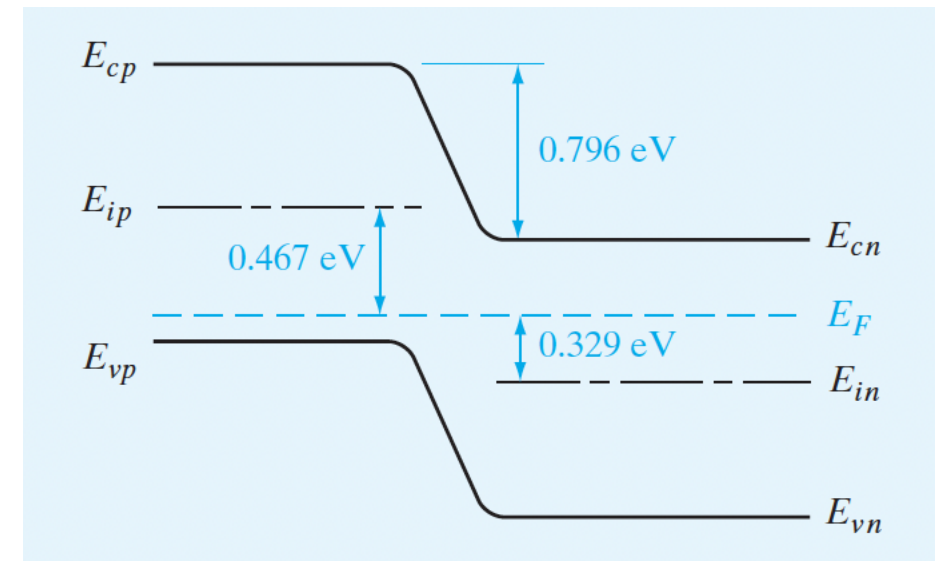
$$E_{ip} - E_F = kT \ln\left(\frac{10^{18}}{1.5 \times 10^{10}}\right) = \underline{\hspace{1cm}} \text{ eV}$$

$$E_F - E_{in} = kT \ln\left(\frac{5 \times 10^{15}}{1.5 \times 10^{10}}\right) = \underline{\hspace{1cm}} \text{ eV}$$

$$qV_0 = (E_{ip} - E_F) + (E_F - E_{in}) = 0.467 + 0.329 = \underline{0.796 \text{ eV}}$$

- Alternatively,

$$\begin{aligned} qV_0 &= kT \ln\left(\frac{N_a N_d}{n_i^2}\right) = 0.026 \ln\left(\frac{(10^{18})(5 \times 10^{15})}{(1.5 \times 10^{10})^2}\right) \\ &= \underline{0.796 \text{ eV}} \end{aligned}$$



# Problem: p-n Junctions in Equilibrium

- Consider the same sample (an abrupt Si p-n junction has  $N_a = 10^{18} \text{ cm}^{-3}$  on one side and  $N_d = 5 \times 10^{15} \text{ cm}^{-3}$  on the other) with a circular cross section of  $10 \text{ } \mu\text{m}$ . Calculate  $x_{n0}$ ,  $x_{p0}$ ,  $Q_+$ , and  $E_0$  for this junction at equilibrium (300 K). Sketch  $E(x)$  and charge density to scale.

$$W = \left[ \frac{2\epsilon V_0}{q} \left( \frac{1}{N_a} + \frac{1}{N_d} \right) \right]^{\frac{1}{2}} = \left[ \frac{(2)(11.8)(8.85 \times 10^{-14})(0.796)}{1.6 \times 10^{-19}} \left( \frac{1}{10^{18}} + \frac{1}{5 \times 10^{15}} \right) \right]^{\frac{1}{2}} = 0.457 \text{ } \mu\text{m}$$

$$x_{n0} = \frac{W}{1 + \frac{N_d}{N_a}} = \frac{0.457}{1 + 5 \times 10^{-3}} = 0.455 \text{ } \mu\text{m}$$

$$x_{p0} = \frac{W}{1 + \frac{N_a}{N_d}} = \frac{0.457}{1 + 200} = 2.27 \times 10^{-3} \text{ } \mu\text{m}$$

$$A = \pi r^2 = \pi(5 \times 10^{-4} \text{ cm})^2 = 7.85 \times 10^{-7} \text{ cm}$$

$$\begin{aligned} Q_+ &= qAx_{n0}N_d \\ &= (1.6 \times 10^{-19})(7.85 \times 10^{-7})(0.455 \times 10^{-4})(5 \times 10^{15}) \\ &= qAx_{p0}N_a = 2.85 \times 10^{-14} \text{ } \end{aligned}$$

$$\begin{aligned} E_0 &= -\frac{qx_{n0}N_d}{\epsilon} = -\frac{q|x_{p0}|N_a}{\epsilon} \\ &= \frac{-(1.6 \times 10^{-19})(0.455 \times 10^{-4})(5 \times 10^{15})}{(11.8)(8.85 \times 10^{-14})} \\ &= 3.48 \times 10^4 \text{ V/cm} \end{aligned}$$

